# Laplace-Spectra as Fingerprints for Image Recognition 

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#### Abstract

In the area of image retrieval from data bases and for copyright protection of large image collections there is a growing demand for unique but easily computable fingerprints for images. These fingerprints can be used to quickly identify every image within a larger set of possibly similar images. This paper introduces a novel method to automatically obtain such fingerprints from an image. It is based on a reinterpretation of an image as a Riemannian manifold. This representation is feasible for gray value images and color images. We discuss the use of the spectrum of eigenvalues of different variants of the Laplace operator as a fingerprint and show the usability of this approach in several use cases. Contrary to existing works in this area we do not only use the discrete Laplacian, but also with a particular emphasis the underlying continuous operator. This allows better results in comparing the resulting spectra and deeper insights in the problems arising. We show how the well known discrete Laplacian is related to the continuous Laplace-Beltrami operator. Furthermore we introduce the new concept of solid height functions to overcome some potential limitations of the method.


## Keywords

Laplace spectra, image recognition, Laplace-Beltrami operator, fingerprints, watermarks, copyright protection, spectra, Laplace-Kirchhoff operator, Riemannian manifolds, color images, isospectrality, features, invariants, image data bases

## 1. INTRODUCTION

One of the main tasks of computer science is to manage large collections of data. In general these collections are maintained using data base management systems. A variety of such systems exists, as well as different strategies for data base management systems to keep track of inserted data. All these strategies require chunks of data to be identified by a unique key, in order to distinguish them from other items stored in the data base. It is common practice to re-

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fer to such identifiers as fingerprints, in analogy to the way a human individual is identified by the prints of its finger tips.

Furthermore with the area of copyright protection a new field of applications has developed recently. In order to identify not licensed copies of protected material there have been efforts to develop watermarks intrinsic to the material in question. In case of shape models such watermarks should be embedded in the geometry of the shape itself and they should be robust against distortions, caused e.g. by reconstructing shapes through scan processes. See [Ko et al., 2003a, Ko et al., 2003b] and [Reuter et al., 2005, Reuter et al., 2006] for recent developments in this area. For images this means that an image that is printed out and then scanned again is still identifiable by its fingerprint in the optimal case. Such fingerprints can be used as watermarks, too.

For some types of data constructing a suitable fingerprint is rather straight forward: E.g. if the data is a collection of english words, it is sufficient to identify each word by its representation as an ASCII- or unicode-string. Well known techniques like tries (cf. [Knuth, 1998]) or the like can then be applied easily. With the data consisting of images there is no such straight forward representation. A number of obstacles arise when trying to construct unique fingerprints for images, e.g.:

- Identical images can be represented in different ways: Even if we require all the images to be given as pixels, color images can be expressed in different color spaces. Thus we would need to restrict ourselves to a fixed color space, i.e. RGB space. For images given in other spaces a conversion needs to be applied.
- Unlike ASCII-strings images may contain minor distortions, e.g. resulting from numerical errors during a color space conversion or compression artifacts. We would like our fingerprints to be robust in the presence of such distortions.
- Pixel images can be given in different resolutions. For some applications it would be desirable that the same image in a different resolution would be identified by the same (or a very similar) fingerprint.
- Sometimes the content of an image is independent of the colors chosen, i.e. color rotation or inversion does
not change the meaning of an image. E.g. the negative of a photograph still represents the same content (although it may look strange). Thus we would need identical fingerprints in those situations.

It is an intricate task to construct fingerprints suitable for image collections. In general it is impossible to take the image itself in its RGB-representation as its fingerprint. Even for a rather tiny image with $200 \times 200$ pixels the resulting fingerprint would be a vector with 40000 entries. There is no efficient way to search a 40000-dimensional space of objects. There has been some effort to overcome these dimensional restrictions recently, however generally speaking the less dimensions are involved, the more efficient a search can be carried out. See [Gaede and Günther, 1998] for a survey of search methods on higher dimensional spaces, see [Berchtold et al., 1996] or [Berchtold et al., 1997] for examples of efficient data structures for such applications. Furthermore there is the possibility to reduce the dimensionality of a feature space afterwards using multi dimensional scaling (MDS), see [Bengio et al., 2004] for an overview on different methods.

In other words we need to find a map from the space of images to a much lower dimensional space without losing information that is relevant for the content. It should be noted, that there are different approaches to automatically construct such a map from a given training set of images, e.g. by MDS. Apart from that, we are looking for fingerprints, that are suitable for different setups without having prior training sets.

There have been different approaches to construct fingerprints, the best known of which are feature vectors. These vectors are constructed from a number of features that can be extracted from an image, including:

- brightness (i.e. mean pixel value) of an image
- contrast (i.e. variance of the pixel values) of an image
- overall roundness of contained shapes
- approximate fractal dimension


## - Fourier transform

- skeletal transforms (e.g. discrete medial axis transforms)

Some of these features do obviously not fulfill our requirements, e.g. changing the brightness or contrast of an image does not change important content of the image. Most of the classical techniques are only suitable for certain special classes of images and not for others. I.e. they cannot be applied in a general situation where there is no further knowledge about the nature of the images available. Of course there is a large number of more elaborate techniques but a detailed discussion would be beyond the scope of this short introduction. See e.g. [Veltkamp and Hagedoorn, 1999] or
[Loncaric, 1998] for an overview of different methods.

In this paper we will introduce a method that works in the general case where there is no restriction concerning the kind of images. We will develop criteria to be met in such a setup, and show that our fingerprints fulfill these criteria. We will explain how to obtain these fingerprints using the spectra of a family of operators known as Laplace or Kirchhoff operators. Only discrete versions of these operators have been used in image processing traditionally, while others (like the Laplace-Beltrami operator) are relatively new in this area. In order to compute the operators' eigenvalues, we will interpret images as Riemannian manifolds. This is a novel approach in the area of image analysis. We will explain the techniques from the point of view of its application on images, although they can be used for more general shapes also. See [Wolter and Friese, 2000] and [Reuter et al., 2005, Reuter et al., 2006] for an introduction to the theory of Laplace spectra in general shape matching.

## 2. RELATED WORK

Using the Laplace operator and more generally using eigenvalues of different operators and matrices derived from this operator is a well known and established technique in the community of shape and image recognition. Most of the applications mentioned in this section use discrete forms of the Laplacian directly, i.e. they are using some kind of admittance matrix.

One of the best known applications of the Laplace spectrum of a graph is graph partitioning. This is useful in areas where one has to find a segmentation of a given mesh, e.g. to identify different components of a scene. See [Mohar, 1997, Chung, 1997] for an overview of the mathematical foundations. Closely related is the application of eigenfunctions to remesh given objects (see [Dong et al., 2005]).

Another popular application of the Laplacian can be found in image processing. A local version of the mesh Laplacian is well established for smoothing of images and meshes, see [Foley et al., 1996, Glassner, 1994, Grady and Schwartz, 2003, Bülow, 2004, Field, 1988, Taubin, 2000] for examples. A variant of the Laplace operator can be used for mesh partitioning and compression. This technique is known as spectral compression, details can be found in [Karni and Gotsman, 2000].

Furthermore the Laplace operator is used for dimensionality reduction of high dimensional data spaces (cf. [Belkin and Niyogi, 2002]). This method could be seen as a complement of the method developed in this paper. Belkin and Niyogi assume the existence of a manifold containing all the objects of a given set of objects represented as points in a feature space. One can assume that these points form a manifold since for most applications the space of possible data depends only on few parameters. E.g. given a set of images shot from an indoor scene by a movable camera the results depend solely on six parameters, i.e. placement and orientation of the camera. Belkin and Niyogi then use some
interpolation technique to form a discrete mesh resembling the assumed manifold containing the objects representations in fetaure space and compute eigenvalues and eigenfunctions of the associated mesh Laplacian. These eigenfunctions can be used to reduce the dimensionality of the embedding feature space to a lower dimension, e.g. six dimensions for the camera example.

The principal difference (aside from using the mesh Laplacian rather than the Laplace-Beltrami operator) is that Belkin and Niyogi compute the eigenvalues of the manifold defined by points in a given feature space, that is each of their points represents an entire manifold in our approach. We use the manifold's eigenvalues as features. The Belkin/Niyogi method is useful when there is a set of features present that has to be reduced for some reasons, while our method is useful when there is a manifold representation for each object possible but no efficient features are known. The Belkin/Niyogi technique can roughly be classified as a multi-dimensional scaling method (MDS) although it is implemented differently from the classical approaches (cf. [Bronstein et al., 2003] and [Bengio et al., 2004]).

Another use of MDS is to map shapes to a canonical signature surface (cf. [Elad (Elbaz) and Kimmel, 2003]). Ideally two isometric shapes should be mapped to an (almost) identical signature surface. It would be possible to use the Belkin/Niyogi technique for this application also.

These uses of the Laplacian show that the application of (non Riemannian!) manifolds in feature space are well established in the image analysis community. On the other hand the duality of the image as a discrete height function and its continuous representation as a surface is also well known in the theory of image processing (see [Glassner, 1995]). Nevertheless there seems to be no transfer between the two fields, i.e. the discussion of an image as a Riemannian manifold. We will see in section 5 how to aquire a Riemannian metric for an image and use it for image classification. This discussion seems to be new.

Using the eigenvalues of a different matrix derived from point sets of two given objects in Euclidean space one can compute a best match of these points. This technique called modal matching was developed by Sclaroff and Pentland [Sclaroff and Pentland, 1995] based on the classical well known technique of momentum matrices. No features are derived from the shapes in this approach instead the chosen points are matched to each other directly.

## 3. FEATURES AND INVARIANTS

In this section we will develop criteria for a fingerprint to be used in image identification. We will make precise what we mean by a feature and an invariant.

Definition 1 Let $A, B$ be sets, with $\operatorname{size}(x)$ being the size of a representation for every $x \in A \cup B$. A map $\iota: A \rightarrow B$ is called a feature map iff $\operatorname{size}(\iota(a))<\operatorname{size}(a)$ for all $a \in A$.

For $x \in \mathbb{R}^{n}$ we simply set $\operatorname{size}(x)=b \cdot n$ for some fixed constant $b$. Definition 1 states, that a feature should take up less space than the object itself.

Definition 2 Let $A, B, I$ be sets, $P=\left(P_{i}: i \in I\right)$ be a partition of $A$ with representatives $p_{i} \in P_{i}, \iota: A \rightarrow B$ be a map (and in our case even a feature map) with

$$
\iota(a)=\iota\left(p_{i}\right) \Leftrightarrow a \in P_{i} \quad \forall i \in I
$$

Then $\iota$ is called a $P$-invariant for $A$.

If $\left|P_{i}\right|=1$ for all $i \in I$ we call $\iota$ a characteristic. There is a relation between partitions and maps: Suppose we have a family of maps $F=\left(f_{j}\right)$ with $f_{j}: A \rightarrow A, j \in J$ for some index set $J$ with $f_{j}$ surjective and $f_{j}^{-1}\left(a_{1}\right) \cap f_{j}^{-1}\left(a_{2}\right)=\emptyset$ for all $a_{1}, a_{2} \in A$. Then there exists a natural partition of $A$ with

$$
a \in P_{i} \Leftrightarrow \exists j \in J: f_{j}(a)=p_{i}
$$

It is therefore perfectly admissible to speak of an $F$-invariant, e.g. for $F$ being the family of isometries of an isometry invariant.

We can now give some criteria for fingerprints feasible for image identification, based on arguments from the introduction:
[FEATURE] Fingerprints should be features, i.e. each fingerprint should have a shorter representation than the associated image.
[ISOMETRY] Fingerprints should be isometry invariants. This is a natural requirement if we are dealing with fingerprints of objects that are metric spaces themselves. Here we will interpret images as Riemannian manifolds (see section 5), this discussion seems to be new in the area of image analysis. In the context of image classification this is motivated by the fact that most content preserving operations are isometries. This includes changes of brightness, rotations, mirror operations, color rotations and inversions.
[SCALING] Fingerprints should optionally be made scaling invariants, in order to identify different resolutions of the same image.

There is another important criterion that is not directly related to invariance:
[SIMILARITY] Fingerprints of similar images should be similar.

This cannot be taken for granted, even small almost isometric changes in the object's geometry may change some feature's values non-continuously. This is especially true for discrete features, e.g. the number of "dark" regions on a photo (the number of components in a level set) and similar features.

Concerning criterion [FEATURE], any method of compression could be used in order to generate fingerprints. This way a compressed representation of an image could be viewed as its fingerprint. Nevertheless, traditional compression schemes tend to change significantly for small changes in the data, thus violating [SIMILARITY]. This behaviour is even wanted in some application, e.g. with MD5 hashes (see [Rivest, 1992]), where small changes to a file are required to cause a significant change in the hash. Nevertheless, methods used especially for geometry compression often change continuously, e.g. if they are based on fourier transformations.

## 4. THE LAPLACE OPERATOR AND THE KIRCHHOFF OPERATOR

In this section we will describe the Laplace operator, also known as Laplacian. First we will take a look at the traditional continuous case, then we will examine the analogous construction in the discrete case.

Definition 3 Given a compact Riemannian manifold $M$ of dimension $n, f: M \rightarrow \mathbb{R}$ with $f \in C^{k}(M)$ and $k \geq 2$. Let grad denote the gradient and div the divergence on the manifold $M$. Then

$$
\Delta f:=\operatorname{div}(\operatorname{grad} f)
$$

defines the Laplace operator $\Delta$ on $M$.

For $M$ being a domain of the Euclidean plane $M \subset \mathbb{R}^{2}$ the Laplace operator reduces to

$$
\begin{equation*}
\Delta f=\left(\frac{\partial^{2} f}{\partial x^{2}}+\frac{\partial^{2} f}{\partial y^{2}}\right) \tag{1}
\end{equation*}
$$

In the non-Euclidean case for a Riemannian metric given by $\left(g_{i j}\right)$ and $n=2$ the Laplace operator can be expressed as (cf. [Chavel, 1984])

$$
\begin{equation*}
\Delta f=\sum_{i=1}^{2} \sum_{j=1}^{2} g^{i j}\left(\frac{\partial^{2} f}{\partial x_{i} \partial x_{j}}-\sum_{k=1}^{2} \Gamma_{i j}^{k} \frac{\partial f}{\partial x_{k}}\right) \tag{2}
\end{equation*}
$$

with $\left(g^{i j}\right):=\left(g_{i j}\right)^{-1}$ and $\Gamma_{i j}^{k}$ being the well known Christoffel symbols (of the second kind, see e.g. [DoCarmo, 1976]):

$$
\Gamma_{i j}^{m}=\frac{1}{2} \sum_{k} g^{k m}\left(\frac{\partial}{\partial_{j}} g_{i k}+\frac{\partial}{\partial_{i}} g_{j k}-\frac{\partial}{\partial_{k}} g_{i j}\right)
$$

Note that these are invariant against isometries and scaling.

One can easily check, that for the Euclidean case equation 2 specializes to equation 1 . We will sometimes refer to the Laplace operator in the non Euclidean case as LaplaceBeltrami operator.

As a special case we could identify the map generating the manifold and the function the operator is applied to, i.e. given coordinate functions $F(u, v):=(x(u, v), y(u, v), z(u, v))$ we calculate $\Delta F$ with respect to the manifold given by $F$. Note that we use a generalized version of the Laplacian here that is defined coordinate wise by $\Delta F:=(\Delta x, \Delta y, \Delta z)$. In
this special case of a "double entry" where $F$ serves the double purpose of representing the manifold and defining a vector field on this manifold we get the well known equation:

$$
\begin{equation*}
\Delta F(u, v)=2 H(u, v) n(u, v) \tag{3}
\end{equation*}
$$

where $H$ is the mean curvature and $n$ is the surface normal of the point $F(u, v)$. This form where the parametrization of the manifold and the argument of the operator are being identified is often used for mesh smoothing (see [Field, 1988, Taubin, 2000] for the discrete case). Equation 3 follows directly from equation 2 and the Gauss equations (see [DoCarmo, 1976]):

$$
\begin{align*}
& \partial_{i} \partial_{j} F=\sum_{k} \Gamma_{i j}^{k} \partial_{k} F+h_{i j} n \\
\Rightarrow \quad & h_{i j} n=\partial_{i} \partial_{j} F-\sum_{k} \Gamma_{i j}^{k} \partial_{k} F \tag{4}
\end{align*}
$$

Here $h_{i j}$ are the coefficients of the second fundamental form. It follows:

$$
\begin{align*}
2 H n & =\sum_{i j} h_{i j} g^{i j} n & & \text { by defini } \\
& =\sum_{i j} g^{i j}\left(\partial_{i} \partial_{j} F-\sum_{k} \Gamma_{i j}^{k} \partial_{k} F\right) & & \text { with (4) } \\
& =\Delta F & & \text { with (2) } \tag{2}
\end{align*}
$$

There is a long tradition in studying the eigenvalues of the Laplace operator. Formally we define an eigenvalue $\lambda$ by:

$$
\begin{equation*}
\Delta f-\lambda f=0 \tag{5}
\end{equation*}
$$

The (multi-)set of possible solutions $\lambda$ to this equation is defined to be the Laplacian spectrum of $M$. This spectrum has a variety of interesting properties some of which make it interesting for image and shape identification.

- The spectrum is an isometry invariant (see [Chavel, 1984]), i.e. if one maps $M$ to an isometric manifold $M^{\prime}$ the spectrum remains unchanged. This fulfills criterion [ISOMETRY]. For experimental results and a detailed discussion we refer to [Reuter et al., 2006].
- Continuous changes of the manifold's geometry result in continuous changes of the spectrum. Furthermore we will demonstrate, that small changes of the geometry yield likewise small changes of the spectrum, see [Courant and Hilbert, 1993], p. 366. This corresponds to criterion [SIMILARITY]. It is important to notice that a topological change of the manifold can change its spectrum radically. I.e. given the full disc, removing an infinitesimal small open disc at the center transforms the surrounding disc into a topological annulus. One will observe a significant change of the spectrum without changing the visual appearance. Luckily this restriction does not apply to images as we will see in section 5 .
- The multiset of eigenvalues form an infinite but countable growing sequence, i.e.

$$
0 \leq \lambda_{1} \leq \lambda_{2} \leq \ldots
$$

Nevertheless a finite number of these eigenvalues is sufficient to distinguish shapes in a practical situation. This is because the smaller eigenvalues correspond to "raw" features of the geometry (like area and boundary length) whereas the higher eigenvalues are related to finer details of the geometry. In a practical setup the number of eigenvalues being sufficient to distinguish images is much smaller than the data needed to present the whole image. This corresponds to criterion [FEATURE].

- For a uniform scaling of the manifold by an factor $a>0$ in every dimension, the spectrum is scaled by $\frac{1}{a^{2}}$. We will show how to use this knowledge to cancel this effect of scaling, thus making the spectrum scaling invariant. This fulfills our criterion [SCALING].

Summarizing the points above we take a finite number of eigenvalues of the Laplace operator as a fingerprint of the given manifold. We will see in section 5 how images are related to manifolds.

### 4.1 The Laplacian with a mass density function

In extension to equation 5 we can introduce a mass density function $\rho$. This function assigns a mass density to each point of the manifold that influences the vibration of the material. Formally we can reformulate the equation as

$$
\begin{equation*}
\Delta f-\lambda \rho f=0 \tag{6}
\end{equation*}
$$

with $\rho: M \rightarrow \mathbb{R}$. A mass density of 1 means the usual density, therefore setting $\rho=1$ everywhere leads to the classical problem in equation 5. The new formulation yields some interesting properties:

- The solutions of the problem depend continuously on changes of the manifold and of the mass density function (see [Courant and Hilbert, 1993], pp. 304).
- Given two isometric manifolds $M$ and $M^{\prime}$ with an isometry $I$ and mass densities $\rho$ and $\rho^{\prime}$ that are conformal to the isometry, that is

$$
I(M)=M^{\prime} \quad \text { and } \quad \rho(m)=\rho^{\prime}(I(m)) \quad \forall m \in M
$$

then the spectra are identical. In different words, for a fixed mass density function the spectrum is isometry invariant.

- The statement above still holds if $I$ is not an isometry but an isospectral transplantation (see section 6.3), i.e. a function that maps a manifold to an isospectral twin.
- Scaling the mass density function by a factor of $1 / k$ results in a spectrum also scaled by $1 / k$. Let $\rho^{\prime}:=k \rho$ and $\Delta f=\lambda \rho f$, then:

$$
\Delta f=\lambda \rho f=\lambda \frac{1}{k} \rho^{\prime} f=\frac{1}{k} \lambda \rho^{\prime} f
$$

For each eigenvalue $\lambda$ of the problem with density $\rho$ an eigenvalue $\lambda^{\prime}:=\frac{1}{k} \lambda$ of the problem with density $\rho^{\prime}$ exists with identical eigenfunction $f$.

- When scaling the manifold by $k$ while keeping the mass density function, one gets

$$
\begin{aligned}
\Delta^{\prime} f & =\frac{1}{k^{2}} \Delta f \text { and } \rho^{\prime}=\rho \\
\Rightarrow \Delta^{\prime} f & =\frac{1}{k^{2}} \Delta f=\frac{1}{k^{2}} \lambda \rho f=\frac{1}{k^{2}} \lambda \rho^{\prime} f
\end{aligned}
$$

Thus the spectrum is scaled by $\frac{1}{k^{2}}$.

### 4.2 Numerical calculations

In order to solve the stated eigenvalue problems 5 and 6 using a computer we need to discretize it. This can be accomplished by transforming the eigenvalue problem into an equivalent variational problem for the manifold $M$ (cf. [Reuter et al., 2006] for the steps involved in the transformation):

$$
\int_{M}\langle\nabla \phi, \nabla f\rangle d \xi=\int_{M} \phi \Delta f d \xi=\lambda \int_{M} \rho \phi f d \xi
$$

Here $\xi$ represents the area element of $M, \phi$ is an arbitrary function from the same space as $f$ and $\rho$ is the mass density function. For the classical problem we set $\rho=1$. Now we approximate $f$ by

$$
f=\sum_{k=1}^{m} c_{k} \phi_{k}
$$

with $\left\{\phi_{1}, \ldots, \phi_{m}\right\}$ being a FEM base for the intended space of solutions. Sometimes we will refer to these functions as form functions. Furthermore we choose $\phi \in\left\{\phi_{1}, \ldots, \phi_{m}\right\}$. This way the problem is transformed into a generalized matrix eigenvalue problem (again see [Reuter et al., 2006] for details):

$$
\begin{align*}
A \vec{c} & =\lambda B \vec{c}  \tag{7}\\
A=\left(a_{j k}\right) & =\left(\int_{G}\left\langle\operatorname{grad} \phi_{j}, \operatorname{grad} \phi_{k}\right\rangle W d u d v\right)  \tag{8}\\
B=\left(b_{j k}\right) & =\left(\int_{G} \rho \phi_{j} \phi_{k} W d u d v\right) \tag{9}
\end{align*}
$$

This problem can be solved using standard numerical libraries.

The existence of boundary conditions can have some impact on the choice of form functions also. Given a Dirichlet boundary condition $f(x)=0$ for all $x \in \Gamma$ with $\Gamma$ being the boundary, given a point $P_{i}$ on the boundary and given that the intended form functions have local character, that is each is associated to a special point and is only defined on a compact surrounding of that point, then the boundary condition mostly implies that the corresponding form function $\phi_{i}$ is zero. This is the case for most sets of popular form functions like piecewise linear function ("hat functions"), polynomial bases etc. In this case every product with $\phi_{i}$ and its derivative is cancelled also. This implies that the corresponding entries in the FEM matrices $A$ and $B$ become zero.

### 4.3 The Laplace-Kirchhoff operator

There is a different version of the Laplace operator in the discrete case. It is defined as follows:

Definition 4 Given a node weighted graph $G=(A, M)$ with $A \in \mathbb{R}^{n \times n}$ being the adjacency matrix of the graph. In case of an (edge-)weighted graph the entry $A_{i j}$ contains the weight of the edge connecting nodes $i$ and $j$. For an unweighted graph the weights are simply 0 or 1 . Most of the time we will use the unweighted case, but there are many applications for the weighted graph also (see e.g. [Dong et al., 2005, Field, 1988, Taubin, 2000]). $M=\operatorname{diag}\left(m_{1}, \ldots, m_{n}\right) \in$ $\mathbb{R}_{+}^{n \times n}$ is the diagonal matrix containing the weights of the nodes, if there are any. Let

$$
D_{i i}=\sum_{j=1}^{n} A_{i j}
$$

define the valence matrix of $G$, i.e. the diagonal matrix containing all degrees of the nodes of $G$. Then the discrete Laplace operator of $G$ is given by

$$
L=L(G):=M^{-1}(D-A)
$$

The discrete Laplace operator is also known from the MatrixTree theorem by G. Kirchhoff and is therefore sometimes referred to as Kirchoff operator. To distinguish the operator from the Laplacian in the continuous case we will refer to it as Laplace-Kirchhoff operator.

The spectrum of the Laplace-Kirchhoff operator is given by the ordinary eigenvalue problem

$$
\begin{equation*}
L x-\lambda x=0 \tag{10}
\end{equation*}
$$

with $x \in \mathbb{R}^{n}$. One can easily check that this is equivalent to

$$
\begin{equation*}
(D-A) x=\lambda M x \Leftrightarrow \operatorname{det}(D-A-\lambda M)=0 \tag{11}
\end{equation*}
$$

which is a generalized eigenvalue problem. The advantage of 11 over the formulation in 10 is, that all matrices are given in symmetric form, whereas $L$ might be non-symmetric in general. This can be easily reformulated to the ordinary eigenvalue problem:

$$
\begin{equation*}
M^{-1 / 2}(D-A) M^{-1 / 2} x=\lambda x \tag{12}
\end{equation*}
$$

For symmetric eigenvalue problems more accurate numerical solutions are available.

The Laplace-Kirchhoff operator can be seen as a special discrete formulation of the Laplace-Beltrami operator. If we choose the FEM base from section 4.2 to be the set $\left\{\phi_{i j}:(i, j) \in \mathbb{Z}^{2}\right\}$ we can define a FEM base for the space of functions defined on $\mathbb{R}^{2}$ with

$$
\begin{aligned}
\phi_{i j}(x, y) & :=\phi_{i}(x) \phi_{j}(y) \\
\phi_{i}(x) & :=\phi(x-i) \\
\phi(x) & := \begin{cases}x+C & \text { for }-1 \leq x<0 \\
-x+C & \text { for } 0 \leq x<1 \\
0 & \text { else }\end{cases}
\end{aligned}
$$



Figure 1: Orthogonal FEM base function
using $C=\frac{3+\sqrt{3}}{6}$ (see figure 1 ). This yields the following matrices (by evaluating equations 8 and 9 , note that we choose pairs of numbers as indices here):

$$
\begin{aligned}
& A=\left(a_{i j, k l}\right)= \begin{cases}4 & \text { for }(i, j)=(k, l) \\
-1 & \text { for }|(i, j)-(k, l)|=1 \\
0 & \text { else }\end{cases} \\
& B=\left(b_{i j, k l}\right)= \begin{cases}\frac{\sqrt{3}-1}{3} & \text { for }(i, j)=(k, l) \\
0 & \text { else }\end{cases}
\end{aligned}
$$

after dividing by a factor of $\frac{\sqrt{3}-1}{3}$ on both sides of the equation. This conforms to the Laplace-Kirchhoff operator of a regular grid, except for a factor of $\frac{\sqrt{3}-1}{3}$ for the $B$ matrix. This factor is of no practical importance, since it only leads to a scaling of the spectrum and can be divided from the spectrum after calculation. Now let there be a mass density function given by

$$
\rho(x, y):=K_{[x][y]}
$$

where $B \in \mathbb{N}^{n \times m}$ is a grey value image. For each square in $\mathbb{R}^{2}$ this function $\rho$ is constant, therefore we get (again after dividing by a factor):

$$
\begin{aligned}
& A=\left(a_{i j, k l}\right)= \begin{cases}4 & \text { for }(i, j)=(k, l) \\
-1 & \text { for }|(i, j)-(k, l)|=1 \\
0 & \text { else }\end{cases} \\
& B=\left(b_{i j, k l}\right)= \begin{cases}K_{i j} \frac{\sqrt{3}-1}{3} & \text { for }(i, j)=(k, l) \\
0 & \text { else }\end{cases}
\end{aligned}
$$

This conforms to the Laplace-Kirchhoff operator with mass density function $\rho$. So far we have constructed discrete operators for the case of a uniformly spaced grid with grid size 1. Now let the grid have a grid size of $1 / w$. We define

$$
\begin{aligned}
& \phi^{[w]}(x):=\phi(w x) \\
& \phi_{i}^{[w]}(x):=\phi^{[w]}(x-i)=\phi(w(x-i))
\end{aligned}
$$

By applying the substitution rule for integration we get:

$$
\begin{aligned}
& \int_{-\infty}^{\infty} \phi_{i}^{[w]} \phi_{j}^{[w]} d x=\frac{1}{w} \int_{-\infty}^{\infty} \phi_{i} \phi_{j} d x \\
& \int_{-\infty}^{\infty} \phi_{i}^{\prime[w]} \phi_{j}^{\prime[w]} d x=\frac{1}{w} \int_{-\infty}^{\infty} w \phi_{i}^{\prime} \phi_{j}^{\prime} d x=\int_{-\infty}^{\infty} \phi_{i}^{\prime} \phi_{j}^{\prime} d x
\end{aligned}
$$

This shows that the mass matrix $B$ is scaled by an additional factor of $1 / w$ in the one dimensional case, for the two dimensional case $B$ is scaled by $1 / w^{2}$ and $A$ is scaled by $1 / w$. Again we can divide by $1 / w$.

Summarizing the above results, the Laplace-Kirchhoff operator can be seen as a discrete approximation of the LaplaceBeltrami operator with mass density operator on a regular grid. Given that the mass matrix $B$ is scaled by $w \frac{\sqrt{3}-1}{3}$ the resulting spectrum is an approximation of the spectrum of the Laplace-Beltrami operator. Note that according to section 4.2 one needs to zero out the matrix entries corresponding to a boundary point to zero if there is a Dirichlet boundary condition.

Instead of interpreting the Laplace-Kirchhoff operator as a special case of the Laplacian resulting from FEM calculations we could also view it as an approximation: Let $B$ be a small geodesic disc with boundary $\partial B$ and center point $p, \phi, \psi: B \rightarrow \mathbb{R}$ with $\psi \equiv 1$ on $B$ and $\phi$ being $C^{2}$-smooth. Then we obtain (cf. [Blaschke and Leichtweiß, 1973])

$$
\begin{aligned}
& \int_{B} \psi \Delta \phi d \xi
\end{aligned}=-\int_{\partial B} \psi \frac{\partial}{\partial n} \phi d s-\int_{B}\langle\nabla \phi, \nabla \psi\rangle d \xi
$$

as $\langle\nabla \phi, \nabla \psi\rangle=\sum_{i j} g^{i j} \partial_{i} \phi \partial_{j} \psi=0$ since $\psi \equiv 1$. This gives:

$$
\begin{equation*}
\int_{B} \Delta \phi d \xi=\int_{\partial B} \frac{\partial}{\partial n} \phi d s \tag{13}
\end{equation*}
$$

It is well known that for the arclength $l$ and area $F$ of a geodesic circle the following approximations hold (with $K(p)$ being the Gaussian curvature at center $p$, cf. [Blaschke and Leichtweiß, 1973], p. 204):

$$
\begin{aligned}
& \lim _{r \rightarrow 0} \frac{2 \pi r-l}{r^{3}}=\frac{\pi}{3} K(p) \\
\Rightarrow & l \approx 2 \pi\left(r-\frac{r^{3}}{6} K(p)\right) \\
& \lim _{r \rightarrow 0} \frac{\pi r^{2}-F}{r^{4}}=\frac{\pi}{12} K(p) \\
\Rightarrow & F \approx \pi\left(r^{2}-\frac{r^{4}}{12} K(p)\right)
\end{aligned}
$$

For a sufficiently small radius of $B \Delta \phi$ is approximately constant, thus we have in geodesic polar coordinates $(r, \alpha)$ :

$$
\begin{align*}
\int_{B} \Delta \phi d \xi & =(\Delta \phi+o(r)) \int_{B} 1 d \xi \\
& =\Delta \phi \underbrace{\left(\pi r^{2}\left(1-\frac{r^{2}}{12} K(p)\right)+o\left(r^{4}\right)\right)}_{\text {Area of the geodesic disk }}+o(r) \\
& =\int_{\partial B} \frac{\partial}{\partial n} \phi d s  \tag{using 13}\\
& =\int_{0}^{2 \pi} \frac{\partial}{\partial n} \phi \underbrace{\sqrt{G(r, \alpha)}}_{\text {Arclength element }} d \alpha \\
& =\int_{0}^{2 \pi} \frac{\partial \phi}{\partial n} \underbrace{\left(\left(r-\frac{r^{3}}{6} K(p)\right)+o\left(r^{4}\right)\right)}_{\text {Arclength element }} d \alpha \\
& =\int_{0}^{2 \pi} \frac{\partial \phi}{\partial n} r+\gamma r^{3} d \alpha \tag{0,1}
\end{align*}
$$

Let $\vec{r}_{\alpha}$ be the cartesian vector of the point with coordinates $(r, \alpha)$. Then we have using a linear approximation:

$$
\begin{aligned}
& \frac{\partial \phi(r, \alpha)}{\partial n} r \\
= & \nabla \phi(r, \alpha) \cdot \vec{r}_{\alpha} \\
= & \nabla \phi(0,0) \vec{r}_{\alpha}+\vec{r}_{\alpha}^{t}\left(\begin{array}{cc}
\phi_{x x}(0,0) & \phi_{x y}(0,0) \\
\phi_{x y}(0,0) & \phi_{y y}(0,0)
\end{array}\right) \vec{r}_{\alpha}+o\left(\left|\vec{r}_{\alpha}\right|^{3}\right) \\
= & \phi(r, \alpha)-\phi(0,0)+o\left(\left|\vec{r}_{\alpha}\right|^{2}\right)
\end{aligned}
$$

Note that here $(x, y)$ denote Riemannian normal coordinates with center $p$ implying that $\Delta \phi(p)=\phi_{x x}(0,0)+\phi_{y y}(0,0)$. Since $\int_{0}^{2 \pi} \nabla \phi(0,0) \vec{r}_{\alpha} d \alpha=0$ and the same is true for the mixed partial derivatives, we get

$$
\begin{aligned}
& \int_{0}^{2 \pi} \frac{\partial \phi(r, \alpha)}{\partial n} r d \alpha \\
= & \int_{0}^{2 \pi} \vec{r}_{\alpha}^{t}\left(\begin{array}{cc}
\phi_{x x}(0,0) & 0 \\
0 & \phi_{y y}(0,0)
\end{array}\right) \vec{r}_{\alpha}+o\left(\left|\vec{r}_{\alpha}\right|^{3}\right) d \alpha \\
= & \int_{0}^{2 \pi} \phi(r, \alpha)-\phi(0,0)+o\left(\left|\vec{r}_{\alpha}\right|^{2}\right) d \alpha
\end{aligned}
$$

Choosing a symmetric discretization of the geodesic disk $B$ with angle steps $\mathcal{D}=\pi / n$ we obtain with $\alpha_{i}:=i \mathcal{D}$

$$
\begin{aligned}
& \left\lvert\, \int_{0}^{2 \pi} \vec{r}^{t}\left(\begin{array}{cc}
\phi_{x x}(0,0) & 0 \\
0 & \phi_{y y}(0,0)
\end{array}\right) \vec{r} d \alpha\right. \\
& \left.-\sum_{i=1}^{2 n} \mathcal{D} \vec{r}_{\alpha_{i}}^{t}\left(\begin{array}{cc}
\phi_{x x}(0,0) & 0 \\
0 & \phi_{y y}(0,0)
\end{array}\right) \vec{r}_{\alpha_{i}} \right\rvert\, \leq r^{2} \frac{2 \pi}{n}\left(\left|\phi_{x x}\right|+\left|\phi_{y y}\right|\right)
\end{aligned}
$$

Using the above equations we get:

$$
\begin{aligned}
& \int_{0}^{2 \pi} \frac{\partial \phi}{\partial n} r+\gamma r^{3} d \alpha \\
= & \sum_{i=1}^{2 n} \mathcal{D}\left(\vec{r}_{\alpha_{i}}^{t}\left(\begin{array}{cc}
\phi_{x x}(0,0) & 0 \\
0 & \phi_{y y}(0,0)
\end{array}\right) \vec{r}_{\alpha_{i}}\right. \\
& \left.+r^{2} \frac{2 \pi}{n}\left(\left|\phi_{x x}\right|+\left|\phi_{y y}\right|\right)+o\left(\left|\vec{r}_{\alpha}\right|^{3}\right)\right) \\
= & \sum_{i=1}^{2 n} \mathcal{D}\left(\phi(i \mathcal{D}, r)-\phi(0,0)+o\left(r^{2}\right)\right)+O\left(\frac{1}{n}\right) r^{2}
\end{aligned}
$$

Therefore we get:

$$
\Delta \phi(p)=\lim _{r \rightarrow 0}\left(\frac{\sum_{i=1}^{2 n} \phi(i \mathcal{D}, r)-\phi(0,0)+O\left(\frac{1}{n}\right) r^{2}+o\left(r^{2}\right)}{\pi r^{2}}\right)
$$

Hence

$$
\Delta \phi(p)=\lim _{r \rightarrow 0}\left(\frac{\sum_{i=1}^{2 n} \phi(i \mathcal{D}, r)-\phi(0,0)}{\pi r^{2}}\right)+O\left(\frac{1}{n}\right)
$$

This corresponds to the well known non-matrix formulation of the Laplace-Kirchhoff operator, i.e. for a function $f=$ $\left(f_{1}, \ldots, f_{n}\right)$ defined on a graph with each $f_{i}$ located at a node $p_{i}$ we get the respective discrete form $\Delta f_{i}=\sum_{j}\left(f_{j}-f_{i}\right) w_{i j}$ for edge weights $w_{i j}$. In matrix formulation this is $\Delta f=$ $-L f$ with $L$ as defined in definition 4.

## 5. IMAGES

In this section we will explain how to represent images in order to compute fingerprints. First we will restrict ourselves to gray value images to keep the description simple, but we will also explain, how to extend the technique developed to images given in arbitrary color spaces.

Definition 5 Let $m, n, g \in \mathbb{N}, G=\{0, \ldots, g-1\}$ and $B \in$ $G^{m \times n}$. Then we call $B$ a discrete gray value image with $g$ steps.

We define a gray value image to be a matrix but as one can easily see, this is the same as defining $B$ to be a discrete height map $B:\{1, \ldots, m\} \times\{1, \ldots, n\} \rightarrow G$. We will use both definitions interchangeable. See figure 2 for an example. We will define a gray value of 0 to represent "black" and $g-1$ to represent "white", with the values in between defining shades of gray. This specification is arbitrary and as we will see later - without effect to the actual calculations.

### 5.1 Fingerprints and the Laplace-Beltrami operator

One way to calculate a fingerprint for an image is to take the image as a discrete height map and transform it into a continuous manifold. Then the eigenvalues of its associated Laplace-Beltrami operator can be used as a fingerprint. At


Figure 2: An image and its height map
first glance it is not clear how to transform the image into a continuous form. An ad hoc technique could be to interpolate the manifold with any kind of 2 D spline, e.g. a tensor product NURBS. However, the interpolation would take some time and there would also be calculation time needed for evaluating the resulting rather large spline representation of the surface.

An obvious solution to calculate the needed coefficients of the first fundamental form is to create a bilinear surface and then use the partial derivatives of that surface. For given pixel values $z_{1}, z_{2}, z_{3}, z_{4}$ the associated patch is:

$$
f(u, v):=(1-v)\left((1-h) z_{1}+h z_{2}\right)+\left(v\left((1-h) z_{3}+h z_{4}\right)\right.
$$

with local coordinates $h, v \in[0,1]$ (see figure 3 ). We get the partial derivatives

$$
\begin{aligned}
& f_{h}=v\left(z_{4}-z_{3}\right)+(1-v)\left(z_{2}-z_{1}\right) \\
& f_{v}=h\left(z_{4}-z_{2}\right)+(1-h)\left(z_{3}-z_{1}\right)
\end{aligned}
$$

and therefore the coefficients

$$
\begin{aligned}
g_{11}= & 1+f_{h}^{2} \\
= & 1+\left(v\left(z_{4}-z_{3}\right)+(1-v)\left(z_{2}-z_{1}\right)\right)^{2} \\
g_{12}= & f_{h} f_{v} \\
= & \left(v\left(z_{4}-z_{3}\right)+(1-v)\right. \\
& \left.\left(z_{2}-z_{1}\right)\right)\left(h\left(z_{4}-z_{2}\right)+(1-h)\left(z_{3}-z_{1}\right)\right) \\
g_{22}= & 1+f_{v}^{2} \\
= & 1+\left(h\left(z_{4}-z_{2}\right)+(1-h)\left(z_{3}-z_{1}\right)\right)^{2}
\end{aligned}
$$

Given this Riemannian metric of an image one can use the finite element method to obtain the eigenvalues. See section 4.2 and [Reuter et al., 2005, Reuter et al., 2006] for the details of this calculations.


Figure 3: Bilinear interpolation of partial derivatives

In section 4 we have noted that a change of the topology of a given object changes its spectrum fundamentally without affecting the point set of the object significantly. For images this is of no relevance since we take each image to be a height function defined on a rectangular area, i.e. the resulting manifold is always homeomorphic to the full disc. Therefore changes in topology cannot occur and we only have to deal with changes in geometry. As noted in section 4 these changes take place continuously, which is the desired behaviour.

Furthermore we have noted that uniform scaling of a manifold by $a$ modifies the associated Laplace-Beltrami spectrum by $1 / a^{2}$. This might be problematic with images since images are usually either scaled along their $x$ and $y$ directions (which represents a uniform scaling of an image) or along their $z$ direction only (which corresponds to a change of contrast). As a result of this, changing e.g. the contrast of an image modifies its spectrum in a rather unpredictable manner. To make the spectrum invariant against contrast changes we need to scale the gray values to the range $\left[0, \ldots, s_{\text {max }}\right]$ where $s_{\text {max }}$ represents the larger value of image width and image height. This way the effect of contrast changes are canceled and a scaling of the image in $x$ and $y$ direction becomes a uniform scaling along all axes. To make this spectrum scaling invariant we could simply divide it by its first non zero eigenvalue. This way scaled spectra become identical. Note that this effect can be achieved also using a different method for computing the similarity of two spectra. E.g. if we use the correlation coefficient instead of a simple Euclidean distance, scaling effects are canceled. See [Reuter et al., 2006] for more detailed insights on this.

### 5.2 Fingerprints from the Laplace-Kirchhoff operator

Another possibility is to choose the eigenvalues of the LaplaceKirchhoff operator as a fingerprint. To accomplish this, the image is transformed into a node weighted graph. Each pixel is interpreted as a node whose weight corresponds to the given gray value. Then two nodes are connected iff they are adjacent in the image. One can choose different models of neighborhoods here, for our experiments we implemented the 4-neighborhood (see figure 4). For the resulting graph the Laplace-Kirchhoff operator according to definition


Figure 4: 4-neighborhood of a pixel

4 is calculated and its eigenvalues are computed numerically. Note that one needs only the $n$ smallest eigenvalues, so efficient techniques like the Lanczos algorithm (see [Cullum and Willoughby, 1985]) can be applied. These eigenvalues possibly normalized to avoid scaling effects are then taken as a fingerprint.

### 5.3 Color images

Usually color images are given in some color space representation. The most commonly used color space is the RGB space, i.e. each color pixel is given by a triple ( $r, g, b$ ) with $r$ representing the red value, $g$ the green value and $b$ the blue value of the pixel. Most color spaces are three dimensional although there are some four dimensional spaces also, e.g. the CMYK space used in printing. We will outline a technique for 3D spaces here, the extension to 4D is straight forward.

Definition 6 Let $m, n, g \in \mathbb{N}, G=\{0, \ldots, g-1\}$ and $B \in$ $(G \times G \times G)^{m \times n}$. Then we call $B$ a discrete RGB value image with $g$ steps.

Just like in the monochrome case we can interpret $B$ as a map $B:\{1, \ldots, m\} \times\{1, \ldots, n\} \rightarrow G \times G \times G$. This is some kind of generalized height function. Like the monochrome image can be understood as a two-manifold embedded in 3D space, we can interpret the color image as a two-manifold embedded in 5D space. The manifold is parametrized by:

$$
(u, v) \mapsto(u, v, r(u, v), g(u, v), b(u, v))
$$

Thus we get the following components for the first fundamental form:

$$
\begin{aligned}
g_{11} & =1+r_{u}^{2}+g_{u}^{2}+b_{u}^{2} \\
g_{22} & =1+r_{v}^{2}+g_{v}^{2}+b_{v}^{2} \\
g_{12}=g_{21} & =r_{u} r_{v}+g_{u} g_{v}+b_{u} b_{v}
\end{aligned}
$$

with $r_{u}$ and $r_{v}$ denoting the partial derivatives of $r(u, v)$ and similar definitions for $g$ and $b$. We can then apply exactly the same interpolation techniques as in the monochrome case in section 5.1 and then use the Laplace Beltrami operator to gain a fingerprint for the color image. See section 6.2 for a discussion of advantages and disadvantages that may arise from this approach.

Interestingly this approach is compatible with the method developed for gray value images in section 5.1. In RGB colour space a gray value image is represented by a colour
image with equal channels $f:=r=g=b$. Thus we get:

$$
\begin{aligned}
g_{11} & =1+3 f_{u}^{2} \\
g_{22} & =1+3 f_{v}^{2} \\
g_{12}=g_{21} & =3 f_{u} f_{v}
\end{aligned}
$$

This can be seen as the calculation done for a gray value image with height function $\sqrt{3} f$. This means that we get the same spectrum as with an ordinary height function $f$ only scaled by $1 / 3$.

## 6. ISOMETRY AND ISOSPECTRALITY

In section 5 we have presented two possibilities for constructing fingerprints of an image applying variants of the Laplace operator. The question we are dealing with now is to what extent a fingerprint is unique for a given image.

### 6.1 Isometry

Let us first discuss the spectrum of the Laplace-Beltrami operator. From section 4 we know, that isometric manifolds share identical spectra. So we would have to find out, in which cases two images can be considered isometric. Furthermore, there are some rare cases where manifolds are isospectral (i.e. they share the same spectrum) but are nonisometric. We will discuss them in section 6.3. Especially, two manifolds generated by the process described in section 5.1 are isometric, if their Riemannian metrics $\left(g_{i j}\right)$ are identical:

Theorem 1 Let the manifolds be defined by height functions $f:[0, m-1] \times[0, n-1] \rightarrow \mathbb{R}$ and $g:[0, m-1] \times[0, n-$ $1] \rightarrow \mathbb{R}$ with $g_{i j}^{f}=g_{i j}^{g}=: g_{i j}$. Then:

$$
f=g+\alpha \text { or } f=-g+\alpha
$$

for any $\alpha \in \mathbb{R}$.

Proof We first show that the partial derivatives of $f$ and $g$ are identical up to their sign:

$$
\begin{aligned}
& g_{11}=1+f_{x}^{2} \quad \wedge \quad g_{11}=1+g_{x}^{2} \\
& g_{22}=1+f_{y}^{2} \quad \wedge \quad g_{22}=1+g_{y}^{2} \\
& g_{12}=g_{21}=f_{x} f_{y} \wedge \quad g_{12}=g_{21}=g_{x} g_{y} \\
& \Rightarrow(1) \quad f_{x}^{2}=g_{x}^{2} \\
& \text { (2) } f_{y}^{2}=g_{y}^{2} \\
& \text { (3) } f_{x} f_{y}=g_{x} g_{y} \\
& \Rightarrow(\text { from } 1) \quad\left|f_{x}\right|=\left|g_{x}\right| \\
& \text { (from 2) } \quad\left|f_{y}\right|=\left|g_{y}\right| \\
& \Rightarrow(\text { with } 3) \quad\left(f_{x}=g_{x} \wedge f_{y}=g_{y}\right) \vee\left(f_{x}=-g_{x} \wedge f_{y}=-g_{y}\right)
\end{aligned}
$$

Now let $f_{x}=g_{x}$ and $f_{y}=g_{y}$. Then we have:

$$
\begin{aligned}
& \text { (1) } f=\int_{x_{1}}^{x_{2}} f_{x} d x+C(y) \\
& \text { (2) } f=\int_{y_{1}}^{y_{2}} f_{y} d y+D(x) \\
& \text { (3) } g=\int_{x_{1}}^{x_{2}} f_{x} d x+E(y) \\
& \text { (4) } g=\int_{y_{1}}^{y_{2}} f_{y} d y+F(x)
\end{aligned}
$$

where $C(y)$ and $E(y)$ are functions depending solely on $y$, $D(x)$ and $F(x)$ solely on $x$. We get:
(5) $0=\int_{x_{1}}^{x_{2}} f_{x} d x-\int_{y_{1}}^{y_{2}} f_{y} d y+C(y)-D(x)$
(6) $0=\int_{x_{1}}^{x_{2}} f_{x} d x-\int_{y_{1}}^{y_{2}} f_{y} d y+E(y)-F(x)$
and finally:
(7) $0=C(y)-D(x)-E(y)+F(x)$

This means that there has to be some constant $\alpha$ fulfilling:

$$
D(x)-F(x)=\alpha=C(y)-E(y)
$$

By substitution in (1) and (3), respectively (2) and (4) we get:

$$
f=g+\alpha
$$

Now let $f_{x}=-g_{x} \wedge f_{y}=-g_{y}$. From an analoguous argument we get:

$$
f=-g+\alpha
$$

There is an alternative proof of theorem 1 using knowledge about Riemannian manifolds: Given two surfaces defined by height functions $f$ and $g$ we can define a solid height function for $f$ by:

$$
S_{f}(x, y ; z):=\left(\begin{array}{c}
x \\
y \\
f(x, y)+z
\end{array}\right) \quad \text { with } z \in \mathbb{R}^{n}
$$

and the same for $g$. One can easily check that its Riemannian (volume) metric is given by:

$$
\left(g_{i j}^{S}\right)=\left(\begin{array}{ccc}
g_{11}^{f} & g_{12}^{f} & f_{x} \\
g_{21}^{f} & g_{22}^{f} & f_{y} \\
f_{x} & f_{y} & 1
\end{array}\right)=\left(\begin{array}{ccc}
1+f_{x}^{2} & f_{x} f_{v} & f_{x} \\
f_{x} f_{v} & 1+f_{v}^{2} & f_{y} \\
f_{x} & f_{y} & 1
\end{array}\right)
$$

We know from our elementary considerations at the beginning of this proof that $\left(f_{x}=g_{x} \wedge f_{y}=g_{y}\right) \vee\left(f_{x}=-g_{x} \wedge f_{y}=\right.$ $-g_{y}$ ), hence for $S_{g}$ we get

$$
\left(g_{i j}^{S}\right)=\left(\begin{array}{ccc}
g_{11}^{g} & g_{12}^{g} & g_{x} \\
g_{21}^{g} & g_{22}^{g} & g_{y} \\
g_{x} & g_{y} & 1
\end{array}\right)
$$




Figure 5: Isometric height functions

This resolves to one of

$$
\left(\begin{array}{ccc}
1+f_{x}^{2} & f_{x} f_{v} & f_{x} \\
f_{x} f_{v} & 1+f_{v}^{2} & f_{y} \\
f_{x} & f_{y} & 1
\end{array}\right) \quad\left(\begin{array}{ccc}
1+f_{x}^{2} & f_{x} f_{v} & -f_{x} \\
f_{x} f_{v} & 1+f_{v}^{2} & -f_{y} \\
-f_{x} & -f_{y} & 1
\end{array}\right)
$$

Therefore $S_{f}$ and $S_{g}$ are isometric. Two solids in Euclidean three space are isometric if and only if they are congruent that is they are related by a series of Euclidean motions (see e.g. [Klingenberg, 1995], p. 88). Hence all their faces including the original height functions must be identical up to Euclidean motions.

Theorem 1 shows, that if two images share the same Riemannian metric, one is either a brighter version of the other, its negative or a combination of both. However in order to be isometric two images do not need to share the same Riemannian metric of their height functions. They could also be rotations or mirror images of each other. This still perfectly conforms to criterion [ISOMETRY]. Furthermore there are some cases where isometries result from a change of parameterization. E.g. let

$$
\begin{aligned}
& f_{1}(x, y):=\sqrt{1-(x-1)^{2}} \\
& f_{2}(x, y):=\frac{x}{2} \sqrt{\pi^{2}-4}
\end{aligned}
$$

One can easily check, that these functions do not have identical Riemannian metrics of their height functions, however their graphs on the parameter space $[0,2] \times[0,2]$ are both isometric to the Euclidean rectangle $[0, \pi] \times[0,2]$. This way they are related by an isometry flattening the cylindrical surface to the rectangle. Figure 5 shows the corresponding manifolds. Another possibility to construct isometric images


Figure 6: Rotated height function


Figure 7: Computation of the rotation angle
is to determine the maximal and minimal partial derivative along one axis and then rotate the image at angles $\alpha$ and $-\alpha$ along the other axis (see figure 6). Since one needs to make sure that the rotated manifold remains describable by a height function, $\alpha$ can be calculated as

$$
\alpha<90-\arctan m_{\max } \text { and } \alpha<90+\arctan m_{\min }
$$

with $m_{\min }$ and $m_{\max }$ being the minimal respectively maximal derivative (see figure 7). Figure 8 shows an image and the same image rotated by $5^{\circ}$. Note that the image needs to be rather blurred to allow even a small possible rotation angle of $5^{\circ}$. For a monochrome image $m_{\text {min }}$ and $m_{\text {max }}$ are determined by the maximal possible difference of two gray values, i.e. $g-1$ :

$$
m_{\min }=-g+1 \text { und } m_{\max }=g-1
$$

(see figure 9 ). Typically $g-1=255$ so it follows:

$$
\alpha<90-89.775312=0.22468818
$$

This gives a maximal rotation angle of $0.22468818^{\circ}$, resulting for an original image width of $w$ in a width of $w \cdot \cos \alpha$ after the rotation. In a typical situation with $g-1=255$ and $\alpha=0.22468818$ this yields a width of $99,999231 \%$ of the original, that is for a total image width of 100000 pixels


Figure 8: Original image and by $5^{\circ}$ rotated version


Figure 9: Maximal gray value difference
less than 1 pixel difference! Effectively the rotation would be invisible.

Therefore in practice one rarely encounters images that meet the above conditions, so isometry usually means that two images are identical up to Euclidean motions and mirror operations.

### 6.2 Isometry and color images

Using the Laplace Beltrami spectrum for color images introduces some differences from the monochrome case that can be both advantageous and disadvantageous. First of all, there are some well known cases of isometric but noncongruent two-manifolds in spaces with more than two dimensions. A popular example is the pair catenoide and helicoide (cf. [DoCarmo, 1976]): These are isometric but not congruent. Since they are isometric we can construct a parametrization so that both share identical $g_{i j}$ values. This way we have defined two color images (namely maps from 2 D to 3 D ) that are isometric but look notably different. This can be stated more formally as follows. Let $f, g: \mathbb{R}^{2} \rightarrow \mathbb{R}^{3}$ be maps from the parameter space to the color space defining the two images. Clearly most of the time the manifolds defined by $f$ and $g$ are degenerate (cf. figure 13), but the above examples show that there are nondegenerate cases also. Suppose the resulting manifolds in $\mathbb{R}^{3}$ are isometric. Then there exists some parametrization so that the $g_{i j}$ values are identical for both manifolds. Suppose $f, g$ are given by such a parametrization. Now define two two-manifolds in 5D by the maps $F(u, v):=(u, v, f(u, v))$ and $G(u, v):=(u, v, g(u, v))$. Then the first fundamental form for the manifold $F$ are given by

$$
g_{11}^{F}=1+g_{11}^{f} \cdot g_{11}^{f} \quad g_{22}^{F}=1+g_{22}^{f} \cdot g_{22}^{f} \quad g_{12}^{F}=g_{12}^{f}
$$

and they are identical for $G$. Figure 10 shows the resulting helicoide and catenoide images with marking lines between the areas where each prime color (red, green and blue) dominates to emphasize the difference. Figure 11 shows the color gamuts of these images. Note that these gamuts approximate the geometry of helicoide and catenoide in RGB-space. Thus we have the disadvantage of non-similar images sharing identical spectra here. Of course one could calculate separate spectra for each of the colors, but in that case we lose the invariance respective color rotations, which we will describe below.

Another possibility to create isometric manifolds associated with color images is, to apply Euclidean motions in color


Figure 10: Color images associated to helicoide and catenoide


Figure 11: Color gamuts of helicoide and catenoide
space, most notably color rotations. In other words: the fingerprints are invariant with respect to color rotations. This can be a big advantage in some applications since color rotations rather represent a change in the look of an image but not in its intended content. Figure 12 shows an example of a color rotation by 90 degrees. Figure 13 shows the two images from figure 12 in RGB-space (their color gamuts). Note that the resulting two-manifolds in 3D are congruent up to rotation. Also note that the gamuts almost fill the entire RGB cube. In such cases color rotation angles are naturally limited to multiples of 90 degrees. Further Euclidean motions in color space correspond to translations. These are effectively linear changes of the dominant image color, once again the principal content of the image remains unchanged. Another advantage is, that the fingerprints are invariant to changes of the color space, given that the color spaces are connected by an Euclidean motion. E.g. one can obtain a representation in CMY space from a RGB representation by


Figure 12: Color rotation


Figure 13: Color gamuts


Figure 14: Two isospectral domains
the transformation

$$
\left(\begin{array}{c}
c \\
m \\
y
\end{array}\right)=\left(\begin{array}{l}
1-r \\
1-g \\
1-b
\end{array}\right)
$$

Clearly this is an Euclidean motion (an inversion and a translation) thus the fingerprint does not change.

### 6.3 Isospectrality

On the other hand there are cases, where two manifolds share an identical spectrum while they are not isometric. We have seen how to represent monochrome images as twomanifolds, thus we have to investigate if it is possible for a pair of two-manifolds to be isospectral. Indeed, examples for non-isometric but isospectral two-manifolds have been found, see [Conway et al., 1994] for some planar domains. Figure 14 shows a typical pair. However, none of those manifolds can be expressed by height functions over a convex domain, i.e. their parameter space must be a non-convex domain. Thus it is very unlikely that there exist isospectral monochrome images, since they are defined by height functions over rectangular (and therefore convex) domains. Figure 15 shows a possible parameter domain for the manifolds from figure 14. From all we know so far isospectrality seems to be a rare phenomenon. Only pairs of isospectral domains have been found in 2D and all of them are nonconvex. Therefore isospectrality should not be a serious problem and should not prevent the discrimination of images via their spectra.

When using the Laplace-Kirchhoff operator instead of the Laplace-Beltrami operator one has to deal with isospectral graphs instead of isospectral manifolds. Just like in the


Figure 15: Parameter domain for the isospectral domains
continuous case the phenomenon is not understood completely but theories exist of how to construct such graphs, see e.g. [Halbeisen and Hungerbühler, 1999]. Halbeisen and Hungerbühler describe a method for obtaining an isospectral "twin" for a given graph. An interesting property of their method is, that the constructed graph is either not simple or has at least two more nodes than the original graph.

## 7. SPECTRA AND FINGERPRINTS

Formally we have defined the Laplacian spectra in equation 5,6 and in equation 10. Following the first definition a spectrum is an infinite countable multiset of real numbers, i.e. a sequence $0 \leq \lambda_{0} \leq \lambda_{1} \leq \ldots \leq \infty$. It is a multiset since one can have eigenvalues with multiplicities larger than 1 . According to the second definition it is a finite multiset of real numbers with $\lambda_{0}=0$. The first eigenvalue is always zero if there is no boundary condition, since the sum of all rows of the matrix $D-A$ is zero, therefore its determinant is zero.

We define the fingerprint of an image to be a suitable finite subset of one of the possible spectra. We will see in section 8 that for practical applications most of the time the first $n \leq 10$ eigenvalues suffice.

To effectively compare two images one has to compare their associated fingerprints. This can be accomplished defining a suitable metric on the space of the fingerprints. If we choose our fingerprints to be the first $n$ eigenvalues the embedding becomes the vector space $\mathbb{R}^{n}$, where we can choose between a large number of well known metrics. For our tests we have used different $p$-norms given by

$$
d_{p}(u, v):=\left(\sum_{i=0}^{n-1}\left|u_{i}-v_{i}\right|^{p}\right)^{\frac{1}{p}}
$$

with $u=\left(u_{1}, \ldots, u_{n}\right)$ and $v=\left(v_{1}, \ldots, v_{n}\right)$ being fingerprints. Furthermore we have tested the Hausdorff distance and the Pearson correlation distance:

$$
d_{c}(u, v):=1-\left|\frac{(n-1) \sum_{i=0}^{n-1}\left(u_{i}-\bar{u}\right)\left(v_{i}-\bar{v}\right)}{\sum_{i=0}^{n-1}\left(u_{i}-\bar{u}\right)^{2} \sum_{i=0}^{n-1}\left(v_{i}-\bar{v}\right)^{2}}\right|
$$

where $\bar{u}$ and $\bar{v}$ denote the arithmetic means of $u$ and $v$. We observed that for most applications the Euclidean distance $d_{2}$ yields acceptable results (see section 8 ) while being easy to implement.


Figure 16: Original test image and modified versions

Since we know that scaling an image can be transformed to an associated scaling of the fingerprint (see sections 4 and 5), one could use a distance that is insensitive against linear scaling. For example the correlation distance has this property by definition. Other possibilities include dividing the fingerprint by its first non zero eigenvalue or normalizing the slope of the best fitting line of the sequence of eigenvalues to 1 before comparing. This way one makes sure that fingerprints being in fact scaled variations of each other are considered identical.

## 8. EXAMPLES

To test the implementation of the methods described above an image from the collection of Snodgrass [Snodgrass and Vanderwart, 1980] was modified (see figure 16). We calculated the eigenvalues of the Laplace-Kirchhoff operator with Dirichlet boundary condition. All images were scaled to $32 \times 32$ pixels to allow a fast computation. The fingerprints were divided by the second eigenvalue (since the first one is always zero), cropped to 10 eigenvalues and compared via Euclidean distance. Figure 17 shows the result for the 15 best matches. The computed distances are shown in the figure, note that the distances plotted are not linearly scaled: the shorter distances are scaled up.

The experiment was repeated using the Laplace-Beltrami operator obtained from the height function. The fingerprints were computed using 338 elements and cubic form functions (cf. [Reuter et al., 2006] for details of the FEM calculations used). They were divided by the first eigenvalue and compared via Euclidean distance. Figure 18 shows the result.

To test the robustness against scaling each image from the collection of Rossion (see [Rossion and Pourtois, 2004]) was scaled by a factor of 2 and added to the collection. The fingerprints were calculated using the Laplace-Beltrami operator obtained from the height function with 338 finite elements. They were compared using Euclidean distances with best fitting lines (see section 7). For 511 of 532 the double sized images were the second best fits (behind the respective image itself), for the remaining 21 images the double sized versions were the third best fits.

The experiment was repeated with the images being changed in contrast instead of size. For each image a copy with $50 \%$ of the original contrast was added to the collection. In this case all of the 532 images could be matched to their low contrast counterparts.


Figure 17: Distances for the Laplace-Kirchhoff operator


Figure 18: Distances for the Laplace-Beltrami operator


Figure 19: Original colour test image and modified versions


Figure 20: Distances for the Laplace-Beltrami operator (colour images)

To test the methods developed for colour images an image from the collection of Kambeck [Kambeck, 2005] was modified in different ways (see figure 19). 30 eigenvalues were calculated using the Laplace-Beltrami operator for generalized height functions with 338 finite elements and cubic form functions. They were compared using the Euclidean distance with best fitting lines. Figure 20 shows the calculated distances for the best matches. Figure 21 shows an MDS plot of the best matches depicted in figure 20. One can clearly identify the cluster of the microscope images. Using the Pearson correlation distance (see section 7) yielded similar good results in this case, whereas using the Hausdorff distance showed inferior results.

Finally we generated a family of images containing a moved full disc (see figure 22). Clearly, this is a one parameter family of images, so one should predict the set of fingerprints to


Figure 21: Cluster for the colour images (LaplaceBeltrami)


Figure 22: One parameter family of images
depend on one parameter as well. Figure 23 shows an MDS plot of the calculated fingerprints. It forms a one parameter curve.

## 9. CONCLUSION AND OUTLOOK

We have presented three methods for obtaining fingerprints from discrete monochrome or color images. Namely the Laplace-Beltrami operator with height functions, the LaplaceBeltrami operator with density functions and its discrete counterpart the Laplace-Kirchhoff operator were used. We have shown an interlink between the discrete Laplace-Kirchhoff operator and the Laplace-Beltrami operator. Further-


Figure 23: One parameter curve of the fingerprints
more we have introduced concepts from the theory of Riemannian geometry into the field of image fingerprints. We have demonstrated these techniques to work in a set of practical situations. The methods were shown to be especially useful in presence of rotations or color rotations of the images in question, since the calculations are invariant against such transformations.

Future work will include investigations on how to apply the Laplace-Kirchhoff operator and the Laplace-Beltrami operator with density functions to color images. Furthermore methods would have to be provided to cancel the effects of special non Euclidean transformations (e.g. most color space transformations) on color images.

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